

I. Χρήσιμες φυσικές σταθερές

Σωματίδιο	Μάζα (Kg)	Φορτίο (Cb)
Ηλεκτρόνιο (e)	$m_e = 9,1091 \times 10^{-31}$	$q_e = -1,6021 \times 10^{-19}$
Πρωτόνιο (p)	$m_p = 1,6725 \times 10^{-27}$	$q_p = +1,6021 \times 10^{-19}$
Νετρόνιο (n)	$m_n = 1,6748 \times 10^{-27}$	0
$ q_e = q_p \approx 1,6 \times 10^{-19} \text{ Cb}$		

$c = 3 \times 10^8 \text{ m/s}$
 $1 \text{ J} = 6,24 \times 10^{18} \text{ eV}$
 $k_B = 1,38 \times 10^{-23} \text{ J/K} = 8,62 \times 10^{-5} \text{ eV/K}$
 $K_e = 9 \times 10^9 \text{ Nm}^2/\text{Cb}^2$
 $h = 6,62 \times 10^{-34} \text{ Js}$
 $\hbar = h/2\pi = 1,05 \times 10^{-34} \text{ Js}$
 $hc = 19,86 \times 10^{-26} \text{ Jm} = 12412 \text{ eVm}$

II. Χρήσιμοι μαθηματικοί τύποι

$$(a + b)^n = a^n + \frac{n}{1!} a^{n-1} b + \frac{n(n-1)}{2!} a^{n-2} b^2 + \dots$$

$$(1 + x)^n = 1 + nx + \frac{n(n-1)}{2!} x^2 + \dots$$

$$e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots$$

$$\ln(1 \pm x) = \pm x - \frac{1}{2}x^2 \pm \frac{1}{3}x^3 - \dots$$

$$\sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \dots$$

$$\cos x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \dots$$

$$\tan x = x + \frac{x^3}{3} + \frac{2x^5}{15} + \dots \quad |x| < \pi/2$$

} x in radians

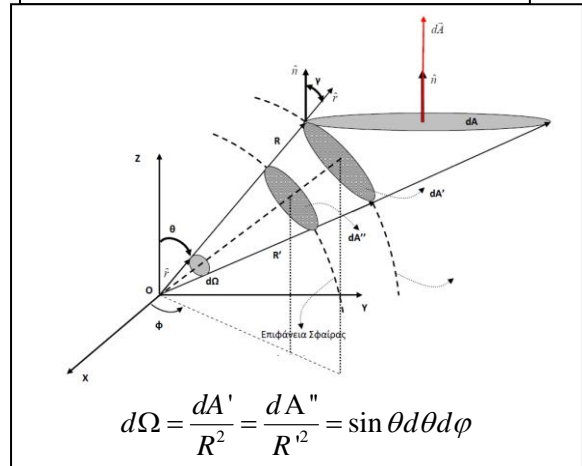
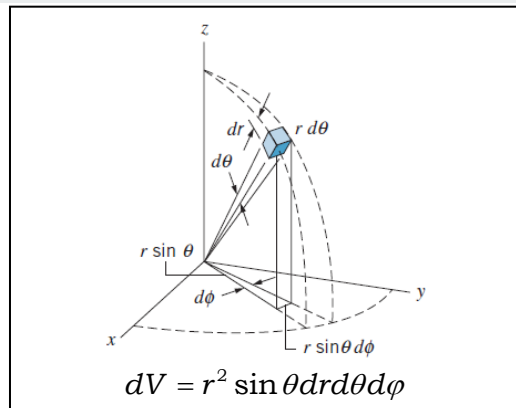
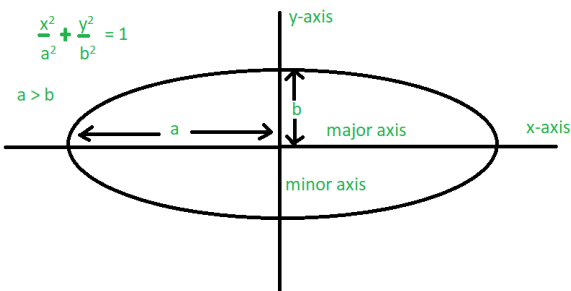
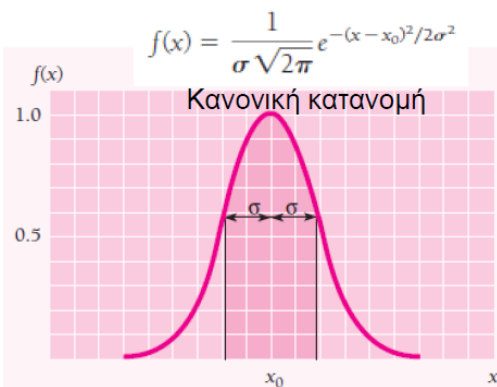
$\sin^2 \theta + \cos^2 \theta = 1$	$\csc^2 \theta = 1 + \cot^2 \theta$
$\sec^2 \theta = 1 + \tan^2 \theta$	$\sin^2 \frac{\theta}{2} = \frac{1}{2}(1 - \cos \theta)$
$\sin 2\theta = 2 \sin \theta \cos \theta$	$\cos^2 \frac{\theta}{2} = \frac{1}{2}(1 + \cos \theta)$
$\cos 2\theta = \cos^2 \theta - \sin^2 \theta$	$1 - \cos \theta = 2 \sin^2 \frac{\theta}{2}$
$\tan 2\theta = \frac{2 \tan \theta}{1 - \tan^2 \theta}$	$\tan \frac{\theta}{2} = \sqrt{\frac{1 - \cos \theta}{1 + \cos \theta}}$
$\sin(A \pm B) = \sin A \cos B \pm \cos A \sin B$	
$\cos(A \pm B) = \cos A \cos B \mp \sin A \sin B$	
$\sin A \pm \sin B = 2 \sin[\frac{1}{2}(A \pm B)] \cos[\frac{1}{2}(A \mp B)]$	
$\cos A + \cos B = 2 \cos[\frac{1}{2}(A + B)] \cos[\frac{1}{2}(A - B)]$	
$\cos A - \cos B = 2 \sin[\frac{1}{2}(A + B)] \sin[\frac{1}{2}(B - A)]$	

Για $x \ll 1$:

$$(1 + x)^n \approx 1 + nx \quad \sin x \approx x$$

$$e^x \approx 1 + x \quad \cos x \approx 1$$

$$\ln(1 \pm x) \approx \pm x \quad \tan x \approx x$$



III. Χρήσιμοι τύποι της Κυματικής

Κυματική εξίσωση: $\frac{\partial^2 y(x,t)}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y(x,t)}{\partial t^2}$

Πυκνότητα ενέργειας κύματος: $u = \frac{\langle E \rangle}{V}$ | Ένταση κύματος: $I = v u$ | $\left(\frac{dE}{dt} \right)_{ave} = IA = v u A$

Κυματοπακέτα:

$$y(x,t) = \sum_i A_i \cos(k_i x - \omega_i t) = R_e \left\{ \sum_i A_i e^{i(k_i x - \omega_i t)} \right\}$$

$$y(x,t) = \int_{k_0 - \Delta k}^{k_0 + \Delta k} A(k) \cos[kx - \omega(k)t] dk = \int_{k_0 - \Delta k}^{k_0 + \Delta k} A(k) e^{i[kx - \omega(k)t]} dk \quad | \quad A(k) = \begin{cases} \neq 0, & k_0 - \Delta k < k < k_0 + \Delta k \\ 0, & \text{αλλού} \end{cases}$$

$$y(x,0) = y_0(x) = \int_{-\infty}^{+\infty} A(k) \cos(kx) dk = \int_{-\infty}^{+\infty} A(k) e^{ikx} dk$$

$$y(x,t) = \int_{-\infty}^{+\infty} A(k) \cos[kx - \omega(k)t] dk = \int_{-\infty}^{+\infty} A(k) e^{i[kx - \omega(k)t]} dk$$

$$v_g = \frac{d\omega(k)}{dk} = \frac{d(v_p k)}{dk} = v_p + k \frac{dv_p}{dk} = v_p - \lambda \frac{dv_p}{d\lambda}$$

$$v_g = \frac{d\omega(k)}{dk} = \left(\frac{d\omega}{dk} \right)_{k=k_0}$$

IV. Χρήσιμοι τύποι της Κλασικής Στατιστικής Φυσικής

$$P = \frac{1}{\zeta} e^{-\frac{\epsilon}{k_B T}} \quad | \quad \zeta = \iiint \iiint e^{-\frac{\epsilon}{k_B T}} dp_x dp_y dp_z dx dy dz \quad | \quad \epsilon = \frac{p^2}{2m} + U(\vec{r})$$

$$\bar{\epsilon} = -\frac{1}{\zeta} \frac{\partial \zeta}{\partial \beta} \quad | \quad \beta = \frac{1}{k_B T}$$

V. Χρήσιμοι τύποι της Θεωρίας Σχετικότητας

Ταχύτητα του φωτός: $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$

Σχετικιστική ορμή: $\vec{p} = \frac{m_0 \vec{v}}{\sqrt{1 - v^2/c^2}}$ | Σχετικιστική μάζα: $m = \frac{m_0}{\sqrt{1 - v^2/c^2}}$

Σχετικιστική κινητική ενέργεια: $K = \frac{m_0 c^2}{\sqrt{1 - v^2/c^2}} - m_0 c^2$

Για ελεύθερο σωματίο: $E = \sqrt{(m_0 c^2)^2 + (cp)^2}$

$$E = K + E_0 = K + m_0 c^2 = \frac{m_0 c^2}{\sqrt{1 - v^2/c^2}} = mc^2$$

$$pc = \sqrt{(K + m_0 c^2)^2 - (m_0 c^2)^2} = \sqrt{K^2 + 2Km_0 c^2} = K \sqrt{1 + \frac{2m_0 c^2}{K}}$$

Για σωματίο σε πεδίο δυνάμεως: $E = K + U + m_0 c^2 = mc^2$

VI. Ηλεκτρομαγνητικά κύματα

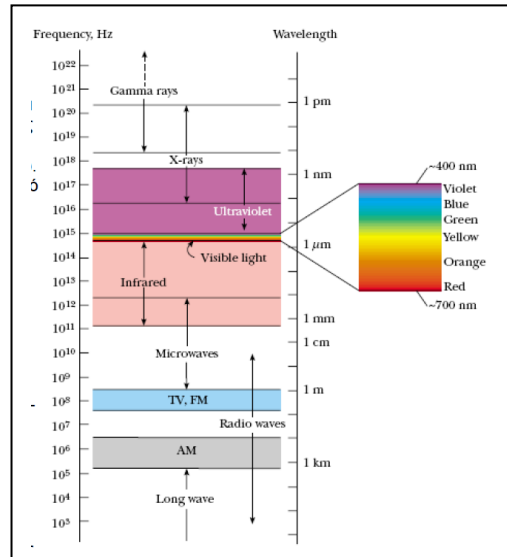
$$\vec{E}(x,t) = E_0 \sin(kx - \omega t) \hat{j}$$

$$\vec{B}(x,t) = B_0 \sin(kx - \omega t) \hat{k}$$

$$\vec{B}(x,t) = \frac{\vec{E}(x,t)}{c}, \quad c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

$$u = \epsilon_0 E^2$$

$$I = cu = c\epsilon_0 E^2$$



VII. Μέλαν σώμα

Αφεικτική Ικανότητα: $\alpha = \frac{P}{A}$, $d\alpha = \rho_T(f)df$, $d\alpha = \rho_T(\lambda)d\lambda$

Πυκνότητα Ενέργειας: $u = \frac{E}{V}$, $du = u_T(f)df$, $du = u_T(\lambda)d\lambda$

$$\rho(f,T) = \frac{c}{4} u(f,T), \quad \rho(\lambda,T) = \frac{c}{4} u(\lambda,T)$$

Νόμος Stefan – Boltzmann: $\alpha = \sigma T^4$, $\sigma = 5,67037 \times 10^{-8} [W / m^2 K^4]$

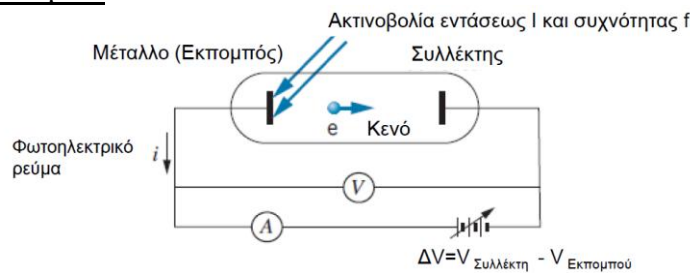
Νόμος Wien: $\lambda_{max} T = \text{σταθερά} = 2,8978 \times 10^{-3} [mK]$, $\frac{f_{max}}{T} = \text{σταθερά} = 5,889 \times 10^{10} [HzK^{-1}]$

Ασυμπτωτικός νόμος Wien: $\rho(f,T) = C_1 f^3 e^{-\frac{C_2 f}{T}}$, για $f \rightarrow \infty$

Νόμος Rayleigh-Jeans: $\rho(f,T) = \frac{2\pi}{c^2} f^2 k_B T$, $\rho(\lambda,T) = \frac{2\pi c}{\lambda^4} k_B T$

Νόμος Planck: $\rho(f,T) = \frac{2\pi h}{c^2} \frac{f^3}{e^{\frac{hf}{k_B T}} - 1}$, $\rho(\lambda,T) = \frac{2\pi hc^2}{\lambda^5} \frac{1}{e^{\frac{hc}{\lambda k_B T}} - 1}$

VIII. Φωτοηλεκτρικό φαινόμενο



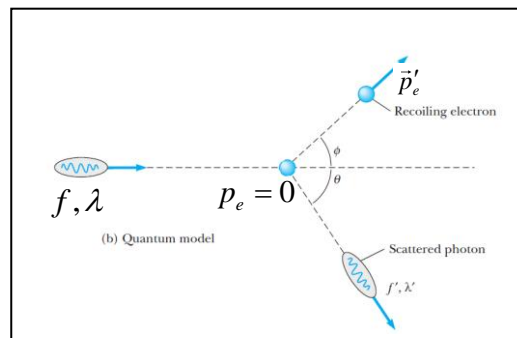
Πειραματική μελέτη: $K_{max} = \frac{1}{2} m_e v_{max}^2 = q_e \Delta V_S$ | Φωτοηλεκτρική εξίσωση του Einstein: $hf = \phi + \frac{1}{2} m_e v^2$

Συχνότητα αποκοπής: $f_c = \frac{\phi}{h}$

IX. Φαινόμενο Compton

$$\Delta\lambda = \lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

Μήκος κύματος Compton: $\lambda_c = \frac{h}{m_e c} \approx 2,4 \times 10^{-2} \text{ \AA}$



X. Φάσμα του Υδρογόνου

$$\frac{1}{\lambda} = R_H \left(\frac{1}{m^2} - \frac{1}{n^2} \right) \left\{ \begin{array}{l} m = 1, 2, 3, 4, \dots \\ n = m + 1, m + 2, m + 3, \dots \end{array} \right.$$

$$f = cR_H \left(\frac{1}{m^2} - \frac{1}{n^2} \right) \left\{ \begin{array}{l} m = 1, 2, 3, 4, \dots \\ n = m + 1, m + 2, m + 3, \dots \end{array} \right.$$

$$R_H = 1,09677576 \times 10^7 \text{ m}^{-1} = \text{σταθερά Rydberg}$$

$$cR_H = 3,29 \times 10^{15} \text{ s}^{-1}$$

XI. Ατομικό πρότυπο του Bohr για τα Μονοηλεκτρονιακά Άτομα

$$L = m_e v r = n \hbar = n \frac{h}{2\pi}, \quad n = 1, 2, 3, \dots$$

$$r_n = \frac{\hbar^2}{Z K_e m_e q_e^2} n^2 = 4\pi\epsilon_0 \frac{\hbar^2}{Z m_e q_e^2} n^2, \quad n = 1, 2, 3, \dots \quad | \quad r_n = \frac{1}{Z} \alpha_0 n^2, \quad n = 1, 2, 3, \dots$$

$$\text{Ακτίνα Bohr: } \alpha_0 = r_1 (Z = 1) = \frac{\hbar^2}{K_e m_e q_e^2} = 4\pi\epsilon_0 \frac{\hbar^2}{m_e q_e^2} \approx 0,53 \times 10^{-8} \text{ cm} = 0,53 \text{ \AA}$$

$$v_n = \frac{Z K_e q_e^2}{\hbar} \frac{1}{n} = \frac{Z q_e^2}{4\pi\epsilon_0 \hbar} \frac{1}{n}, \quad n = 1, 2, 3, \dots$$

$$E_n = -\frac{Z^2 K_e^2 m_e q_e^4}{2\hbar^2 n^2} = -\frac{Z^2 K_e^2 2\pi^2 m_e q_e^4}{h^2 n^2} = -\frac{Z^2 m_e q_e^4}{8\epsilon_0^2 h^2 n^2} = -\frac{13,6}{n^2} Z^2 \text{ eV}, \quad n = 1, 2, 3, \dots$$

$$R_\infty = \frac{m_e q_e^4}{8c\epsilon_0^2 h^3} = 1,0973732 \times 10^7 \text{ m}^{-1}$$

$$\text{Σχετικιστική διόρθωση - Σταθερά λεπτής υφής: } \alpha = \frac{K_e q_e^2}{\hbar c} = \frac{1}{137}, \quad v_n = \frac{Z}{n} \frac{c}{137}$$

XII. Κανόνας κβάντωσης Wilson-Sommerfeld

$$\oint p_q dq = nh, \quad n = 0, 1, 2, 3, \dots$$

XIII. Σύγχρονη Κβαντομηχανική

$$i\hbar \frac{\partial \Psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x,t)}{\partial x^2} + U(x)\Psi(x,t)$$

$$i\hbar \frac{\partial \Psi(\vec{r},t)}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r},t) + U(\vec{r})\Psi(\vec{r},t), \quad \nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Για σωματίο καθορισμένης ενέργειας:

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \Psi(x,t)}{\partial x^2} + U(x)\Psi(x,t) = E\Psi(x,t)$$

$$-\frac{\hbar^2}{2m} \nabla^2 \Psi(\vec{r},t) + U(x)\Psi(\vec{r},t) = E\Psi(\vec{r},t)$$

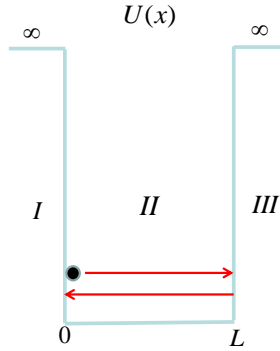
$$\Psi(x,t) = \psi(x) e^{-\frac{i}{\hbar} Et} \quad | \quad \frac{d^2 \psi(x)}{dx^2} + \frac{2m}{\hbar^2} [E - U(x)] \psi(x) = 0$$

XIV. Απειρόβαθο Πηγάδι

$$U(x) = \begin{cases} 0, & 0 \leq x \leq L \\ \infty, & \text{αλλού} \end{cases}$$

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L}x\right), n = 1, 2, 3, \dots$$

$$E_n = n^2 \frac{\pi^2 \hbar^2}{8mL^2}, n = 1, 2, 3, \dots$$



XVI. Αρμονικός Ταλαντωτής

$$U(x) = \frac{1}{2} \kappa x^2 = \frac{1}{2} m \omega^2 x^2, \quad \omega = \sqrt{\frac{\kappa}{m}}$$

$$\psi_n(x) = \left(\frac{\alpha}{\pi}\right)^{\frac{1}{4}} \left(\frac{2^{-n}}{n!}\right)^{\frac{1}{2}} e^{-\frac{\alpha x^2}{2}} H_n\left(\alpha^{\frac{1}{2}}x\right), n = 0, 1, 2, 3, \dots$$

$$H_n(y) = (-1)^n e^{y^2} \frac{d^n (e^{-y^2})}{dy^n}$$

$$\alpha = \frac{m\omega}{\hbar}$$

$$E = \left(n + \frac{1}{2}\right) \hbar \omega, n = 0, 1, 2, 3, \dots$$

n	H _n (y)
0	1
1	2y
2	4y ² - 2
3	8y ³ - 12y
4	16y ⁴ - 48y ² + 12
5	32y ⁵ - 160y ³ + 120y

XVII. Μονοηλεκτρονιακά Άτομα στα πλαίσια της Σύγχρονης Κβαντομηχανικής

1. Κυματοσυναρτήσεις:

$$\psi_{n\ell m_\ell}(\mathbf{r}, \theta, \varphi) = R_{n\ell}(r) \Theta_{\ell m_\ell}(\theta) \Phi_{m_\ell}(\varphi) = R_{n\ell}(r) Y_{\ell m_\ell}(\theta, \varphi)$$

$$\frac{1}{r^2} \frac{d}{dr} \left(r^2 \frac{dR}{dr} \right) + \left[\frac{2m_e}{\hbar^2} \left(E + \frac{1}{4\pi\epsilon_0} \frac{Zq_e^2}{r} \right) - \frac{\ell(\ell+1)}{r^2} \right] R = 0$$

$$\frac{1}{\sin\theta} \frac{d}{d\theta} \left(\sin\theta \frac{d\Theta}{d\theta} \right) + \left[\ell(\ell+1) - \frac{m_\ell^2}{\sin^2\theta} \right] \Theta = 0$$

$$\frac{d^2\Phi}{d\varphi^2} + m_\ell^2 \Phi = 0$$

Πυκνότητες πιθανότητας:

$$P(r) = |R_{n\ell}(r)|^2 r^2, \quad P(\theta) = |\Theta_{\ell m_\ell}(\theta)|^2, \quad P(\varphi) = |\Phi_{m_\ell}(\varphi)|^2, \quad |Y_{\ell m_\ell}(\theta, \varphi)|^2 = \frac{1}{2\pi} |\Theta_{\ell m_\ell}(\theta)|^2$$

2. Στροφορμή:

$$L = \sqrt{\ell(\ell+1)} \hbar \quad L_z = m_\ell \hbar$$

$$\ell = 0, 1, 2, 3, \dots, (n-1)$$

$$m_\ell = -\ell, \dots, 0, \dots, \ell$$

$$\tilde{\mu}_L = -\frac{|q_e|}{2m_e} \vec{L} \quad | \quad \mu_{L,z} = -\frac{|q_e|}{2m_e} L_z = -\frac{|q_e|}{2m_e} m_\ell \hbar = -m_\ell \mu_B \quad | \quad \mu_B = \frac{|q_e| \hbar}{2m_e} = 9,274 \times 10^{-24} \text{ J / T}$$

3. Spin:

$$S = \sqrt{s(s+1)}\hbar, \quad s = \frac{1}{2}$$

$$S_z = m_s \hbar$$

$$m_s = \pm s = \pm \frac{1}{2}$$

$$\vec{\mu}_S = -\frac{|q_e|}{m_e} \vec{S} \quad | \quad \mu_{S,z} = -2\mu_B m_s = \pm \mu_B$$

4. Σύνθεση Στροφορμών:

$$\vec{J} = \vec{L} + \vec{S}$$

$$J = \sqrt{j(j+1)}\hbar \quad J_z = m_j \hbar$$

$$j = \ell + s, \dots, |\ell - s|$$

$$m_j = -j, \dots, 0, \dots, j$$

$$\vec{\mu} = \vec{\mu}_L + \vec{\mu}_S \Rightarrow \vec{\mu} = -\frac{|q_e|}{2m_e} (\vec{L} + 2\vec{S})$$

5. Επιτρεπόμενες Μεταβάσεις: $\Delta \ell = \pm 1, \Delta m_\ell = 0, \pm 1$

ΑΤΟΜΟ ΥΔΡΟΓΟΝΟΥ

n	l	m_l	$\Phi(\phi)$	$\Theta(\theta)$	$R(r)$	$\psi(r, \theta, \phi)$
1	0	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{1}{\sqrt{2}}$	$\frac{2}{a_0^{3/2}} e^{-r/a_0}$	$\frac{1}{\sqrt{\pi} a_0^{3/2}} e^{-r/a_0}$
2	0	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{1}{\sqrt{2}}$	$\frac{1}{2\sqrt{2} a_0^{3/2}} \left(2 - \frac{r}{a_0}\right) e^{-r/2a_0}$	$\frac{1}{4\sqrt{2\pi} a_0^{3/2}} \left(2 - \frac{r}{a_0}\right) e^{-r/2a_0}$
2	1	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{\sqrt{6}}{2} \cos \theta$	$\frac{1}{2\sqrt{6} a_0^{3/2}} \frac{r}{a_0} e^{-r/2a_0}$	$\frac{1}{4\sqrt{2\pi} a_0^{3/2}} \frac{r}{a_0} e^{-r/2a_0} \cos \theta$
2	1	± 1	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$	$\frac{\sqrt{3}}{2} \sin \theta$	$\frac{1}{2\sqrt{6} a_0^{3/2}} \frac{r}{a_0} e^{-r/2a_0}$	$\frac{1}{8\sqrt{\pi} a_0^{3/2}} \frac{r}{a_0} e^{-r/2a_0} \sin \theta e^{\pm i\phi}$
3	0	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{1}{\sqrt{2}}$	$\frac{2}{81\sqrt{3} a_0^{3/2}} \left(27 - 18\frac{r}{a_0} + 2\frac{r^2}{a_0^2}\right) e^{-r/3a_0}$	$\frac{1}{81\sqrt{3\pi} a_0^{3/2}} \left(27 - 18\frac{r}{a_0} + 2\frac{r^2}{a_0^2}\right) e^{-r/3a_0}$
3	1	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{\sqrt{6}}{2} \cos \theta$	$\frac{4}{81\sqrt{6} a_0^{3/2}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0}$	$\frac{\sqrt{2}}{81\sqrt{\pi} a_0^{3/2}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0} \cos \theta$
3	1	± 1	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$	$\frac{\sqrt{3}}{2} \sin \theta$	$\frac{4}{81\sqrt{6} a_0^{3/2}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0}$	$\frac{1}{81\sqrt{\pi} a_0^{3/2}} \left(6 - \frac{r}{a_0}\right) \frac{r}{a_0} e^{-r/3a_0} \sin \theta e^{\pm i\phi}$
3	2	0	$\frac{1}{\sqrt{2\pi}}$	$\frac{\sqrt{10}}{4} (3 \cos^2 \theta - 1)$	$\frac{4}{81\sqrt{30} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\frac{1}{81\sqrt{6\pi} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0} (3 \cos^2 \theta - 1)$
3	2	± 1	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$	$\frac{\sqrt{15}}{2} \sin \theta \cos \theta$	$\frac{4}{81\sqrt{30} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\frac{1}{81\sqrt{\pi} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0} \sin \theta \cos \theta e^{\pm i\phi}$
3	2	± 2	$\frac{1}{\sqrt{2\pi}} e^{\pm 2i\phi}$	$\frac{\sqrt{15}}{4} \sin^2 \theta$	$\frac{4}{81\sqrt{30} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\frac{1}{162\sqrt{\pi} a_0^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0} \sin^2 \theta e^{\pm 2i\phi}$