<u>ΝΑΝΟΗΛΕΚΤΡΟΝΙΚΗ & ΚΒΑΝΤΙΚΕΣ ΠΥΛΕΣ</u> 2^η Διάλεξη

<u>Βιβλιογραφία</u>: EXPLORATIONS IN QUANTUM COMPUTING, Colin P. Williams (2nd edition, Springer-Verlag, 2011), chapter 1.

Quantization: From Bits to Qubits

- Ket Vector Representation of a Qubit
- Superposition States of a Single Qubit
- Bloch Sphere Picture of a Qubit
- Reading the Bit Value of a Qubit

Assumptions about the properties of bit that are no longer necessarily true at the quantum scale

| Assumption | Classically | Quantum mechanically |
|--|-------------|---|
| A bit always has a definite value | True | False. A bit need not have a definite value until the moment after it is read |
| A bit can only be 0 or 1 | True | False. A bit can be in a superposition of 0 and 1 simultaneously |
| A bit can be copied without affecting its value | True | False. A qubit in an unknown state cannot be copied without necessarily changing its quantum state |
| A bit can be read without affecting its value | True | False. Reading a qubit that is initially in a superposition will change the qubit |
| Reading one bit in the computer memory has no affect on any other (unread) bit in the memory | True | False. If the bit being read is entangled with another qubit, reading one qubit will affect the other |
| To compute the result of a computation, you must run the computer | True | False |

Quantization: From Bits to Qubits

Ket Vector Representation of a Qubit: $|0\rangle \equiv \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |1\rangle \equiv \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

A quantum system can be found to be in one of a discrete set of states: $|0\rangle$ or $|1\rangle$

Superposition

If it is not observed it may also exist in a superposition of those states simultaneously: $|\psi\rangle = a|0\rangle + b|1\rangle$ such that $|a|^2 + |b|^2 = 1$.

Dirac Notation:

For every "ket" $|\psi\rangle$ (column vector), there is a corresponding "bra" $\langle\psi|$ (row vector):

$$\begin{aligned} |\psi\rangle &= a|0\rangle + b|1\rangle = \begin{pmatrix} a \\ b \end{pmatrix} \\ \langle\psi| &= a^*\langle 0| + b^*\langle 1| = (a^* \ b^*) \end{aligned}$$

The ket and the bra contain *equivalent information* about the quantum state

Inner and Outer Products

For a pair of qubits in states: $|\psi\rangle = a|0\rangle + b|1\rangle$ and $|\phi\rangle = c|0\rangle + d|1\rangle$

The *inner product* $\langle \psi | \phi \rangle$ defines the **overlap** between (normalized) states: $|\psi\rangle |\phi\rangle$

$$\langle \psi | \phi \rangle = \underbrace{(\langle \psi |) \cdot (| \phi \rangle)}_{\text{bra (c) ket}} = (a^* \ b^*) \cdot \begin{pmatrix} c \\ d \end{pmatrix} = a^*c + b^*d$$
$$\langle \psi | \psi \rangle = (a^* \ b^*) \cdot \begin{pmatrix} a \\ b \end{pmatrix} = a^*a + b^*b = |a|^2 + |b|^2 = 1$$

The outer product $|\psi\rangle\langle\phi|$ is a matrix:

$$|\psi\rangle\langle\phi| = (|\psi\rangle) \cdot (\langle\phi|) = \begin{pmatrix}a\\b\end{pmatrix} \cdot (c^* \ d^*) = \begin{pmatrix}ac^* \ ad^*\\bc^* \ bd^*\end{pmatrix}$$

The outer product describes the structure of **unitary operators**, which correspond to **quantum logic gates**. For example, a NOT gate:

$$NOT = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = |0\rangle\langle 1| + |1\rangle\langle 0|$$

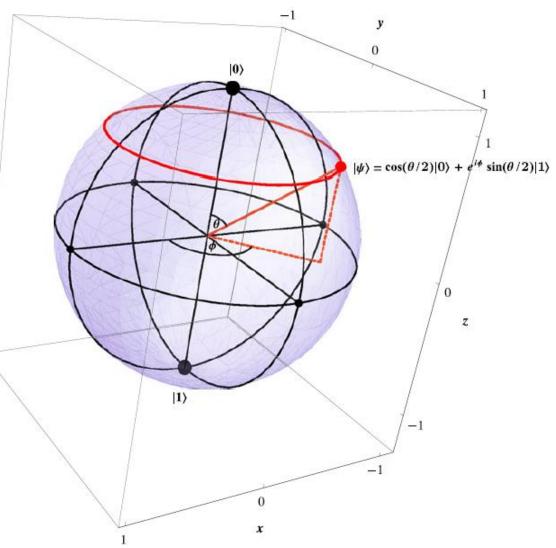
Bloch Sphere Picture of a Qubit

A pure quantum state of a single qubit is a unit vector in Bloch sphere.

A pair of elevation and azimuth angles (θ, ϕ) in the range $0 \le \theta \le \pi$ and $0 \le \phi \le 2\pi$ pick out a point on the Bloch sphere.

Orthogonal states, $|0\rangle$ and $|1\rangle$, are not found to be at right angles on the Bloch sphere.

Orthogonal quantum states, i.e. states $|\psi\rangle$ and $|\chi\rangle$ for which $\langle\psi|\chi\rangle$ = 0, are represented by antipodal points on the Bloch sphere (rather than being drawn at right angles).



Bloch sphere showing the computational basis states $|0\rangle$ and $|1\rangle$, and a general qubit state: $|\psi\rangle = \cos(\theta/2)|0\rangle + e^{i\phi}\sin(\theta/2)|1\rangle$, where θ , and ϕ are real numbers.

Pure 1-qubit states on Bloch Sphere

Bloch sphere labeled with pure 1qubit states at the extremes of the x-, y-, and z-axes:

X-axis: $|\nearrow\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ $|\swarrow\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$

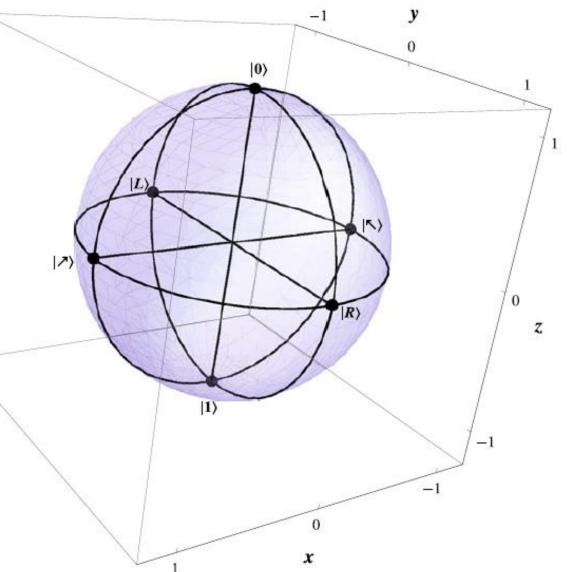
Y-axis:

$$\begin{split} |R\rangle &= |\circlearrowright\rangle = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle) \\ |L\rangle &= |\circlearrowright\rangle = \frac{1}{\sqrt{2}}(|0\rangle - i|1\rangle) \end{split}$$

Z-axis: $|0\rangle$, and $|1\rangle$

Orthogonal quantum states are located at antipodal points.

The operation that maps an unknown state to its antipodal state **cannot be expressed as a rotation** on the Bloch sphere.



Rather it is the sum of a rotation (in longitude through 180 degrees) and a reflection (in latitude with respect to the equatorial plane of the Bloch sphere). 6

Reading the Bit Value of a Qubit

Measuring the bit value of a qubit initially in state: $a|0\rangle + b|1\rangle$ yields the answer:

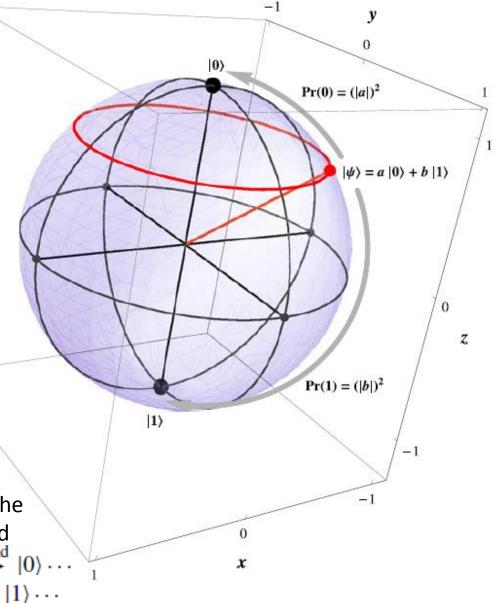
 $\operatorname{Read}(a|0\rangle + b|1\rangle) = \begin{cases} 0 & \text{with probability } |a|^2\\ 1 & \text{with probability } |b|^2 \end{cases}$

and projects the qubit into either state $|0\rangle$ or state $|1\rangle$ respectively.

A measurement of a qubit with respect to North and South poles axis is called a **measurement "in the computational basis"** because the answer we get will be one of the bit values $|0\rangle$ or $|1\rangle$.

Thus, for a single qubit quantum memory register the outcome we obtain from reading it is non-deterministic.

Quantum Zeno Effect: measuring repeatedly the same state its evolution can be suppressed and fixed in a quantum state, $|\psi\rangle \xrightarrow{\text{read}} |0\rangle \xrightarrow{\text{read}} |0\rangle \xrightarrow{\text{read}} |0\rangle \xrightarrow{\text{read}} |0\rangle \cdots$ or $|\psi\rangle \xrightarrow{\text{read}} |1\rangle \xrightarrow{\text{read}} |1\rangle \xrightarrow{\text{read}} |1\rangle \cdots$



7

Problems with Solutions

<u>Problem 1</u> Consider the three cases:

(i)
$$|0\rangle := \begin{pmatrix} 1\\0 \end{pmatrix}, \quad |1\rangle := \begin{pmatrix} 0\\1 \end{pmatrix}$$

(ii) $|0\rangle := \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\1 \end{pmatrix}, \quad |1\rangle := \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\-1 \end{pmatrix}$
(iii) $|0\rangle := \begin{pmatrix} \cos\theta\\\sin\theta \end{pmatrix}, \quad |1\rangle := \begin{pmatrix} \sin\theta\\-\cos\theta \end{pmatrix}$

Find the matrix representation of **A** in these bases. $A := |0\rangle\langle 0| + |1\rangle\langle 1|$

Solution 1. We find:

(i)
$$A = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

(ii)
$$A = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

(iii)
$$A = \begin{pmatrix} \cos^2 \theta & \cos \theta \sin \theta \\ \cos \theta \sin \theta & \sin^2 \theta \end{pmatrix} + \begin{pmatrix} \sin^2 \theta & -\cos \theta \sin \theta \\ -\cos \theta \sin \theta & \cos^2 \theta \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

For all three cases: $A = I_2$ where I_2 is the 2 x 2 unit matrix.

The third case contains the first two as special cases.

Problems with Solutions

<u>Problem 2</u> The NOT operation (unitary operator) is defined as: $|0\rangle \rightarrow |1\rangle$, $|1\rangle \rightarrow |0\rangle$

(i) Find the unitary operator U_{NOT} which implements the NOT operation with respect to the basis $\{ |0\rangle, |1\rangle \}$. (ii) Let

$$|0
angle = \left(egin{array}{c} 1 \ 0 \end{array}
ight), \qquad |1
angle = \left(egin{array}{c} 0 \ 1 \end{array}
ight)$$

Find the matrix representation of U_{NOT} for this basis. (iii) Let

$$|0\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\1 \end{pmatrix}, \qquad |1\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1\\-1 \end{pmatrix}.$$

Find the matrix representation of U_{NOT} for this basis.

Solution 2.

(i) Obviously, $U_{NOT} = |0\rangle\langle 1| + |1\rangle\langle 0|$

since $\langle 0|0\rangle = \langle 1|1\rangle = 1$ and $\langle 0|1\rangle = \langle 1|0\rangle = 0$.

(ii) For the standard basis we find

Thus, the respective matrix representations for the two bases are different.

(iii) For the Hadamard basis we find

$$U_{NOT} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} . \qquad \qquad U_{NOT} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Problems with Solutions

<u>Problem 3</u> The qubit *trine* is defined by the following states:

$$|\psi_0
angle = |0
angle, \quad |\psi_1
angle = -rac{1}{2}|0
angle - rac{\sqrt{3}}{2}|1
angle, \quad |\psi_2
angle = -rac{1}{2}|0
angle + rac{\sqrt{3}}{2}|1
angle$$

where $\{ |0\rangle, |1\rangle \}$ is an orthonormal basis set. Find

 $|\langle \psi_0 | \psi_1 \rangle|^2$, $|\langle \psi_1 | \psi_2 \rangle|^2$, $|\langle \psi_2 | \psi_0 \rangle|^2$.

Solution 3.

Using
$$\langle 0|0\rangle = 1$$
, $\langle 1|1\rangle = 1$ and $\langle 0|1\rangle = 0$ we find
 $|\langle \psi_0|\psi_1\rangle|^2 = \frac{1}{4}$, $|\langle \psi_1|\psi_2\rangle|^2 = \frac{1}{4}$, $|\langle \psi_2|\psi_0\rangle|^2 = \frac{1}{4}$

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